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Mode I crack morphology and fracture surface roughness of granite under impact

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Abstract

To study the response of mode I cracks under impact loading, the dynamic fracture test on the single cleavage triangle (SCT) granite specimens with the lateral opening was carried out using a split Hopkinson pressure bar test system. The Hough transform method was used to quantitatively describe the distributions of the surface crack lengths and angles, and the relationship between them and the absorbed energy of the specimen was analyzed. The three-dimensional (3D) point cloud data of the fracture surface was obtained by using a 3D surface topography detection device, and a fracture surface reconstruction method based on the threshold detection of the fitting surface was proposed, which effectively solved the error problem caused by the appearance of error points in the 3D point cloud data of the fracture surface. The relationship between fracture surface roughness and absorbed energy was also discussed. The results show that the exponential distribution of crack length λ tends to increase, and the distribution of crack angles is more even with the increase in absorbed energy. There are relatively few cracks in the horizontal direction, and the cracks have obvious directionality when the dissipated energy of the specimen is lower, while the bending degree of the cracks is higher and the connectivity is better when the dissipated energy is higher. The fracture surface reconstruction method can complete the fracture surface processing better at 0.03-0.08 of the points in the x and y directions and the threshold of 0.25. The roughness statistical parameters of fracture surfaces A and B showed a decreasing trend with the increase in dissipated energy.

Keywords

single cleavage triangle (SCT) specimen, rock fracturing, Hough transform, fracture surface roughness, mode I crack

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Mode I crack morphology and fracture surface roughness of granite under impact

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Abstract: To study the response of mode I cracks under impact loading, the dynamic fracture test on the single cleavage triangle (SCT) granite specimens with the lateral opening was carried out using a split Hopkinson pressure bar test system. The Hough transform method was used to quantitatively describe the distributions of the surface crack lengths and angles, and the relationship between them and the absorbed energy of the specimen was analyzed. The three-dimensional (3D) point cloud data of the fracture surface was obtained by using a 3D surface topography detection device, and a fracture surface reconstruction method based on the threshold detection of the fitting surface was proposed, which effectively solved the error problem caused by the appearance of error points in the 3D point cloud data of the fracture surface. The relationship between fracture surface roughness and absorbed energy was also discussed. The results show that the exponential distribution of crack length λ tends to increase, and the distribution of crack angles is more even with the increase in absorbed energy of the specimen is lower, while the bending degree of the cracks is higher and the connectivity is better when the dissipated energy is higher. The fracture surface reconstruction method can complete the fracture surface processing better at 0.03-0.08 of the points in the *x* and *y* directions and the threshold of 0.25. The roughness statistical parameters of fracture surfaces A and B showed a decreasing trend with the increase in dissipated energy.

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1 Introduction

In geotechnical engineering such as mine, water conservancy and tunnel, dynamic fracture of rock mass is often triggered by blasting and other dynamic loads, resulting in instability and failure of underground structures^[1]. As a heterogeneous material, rock contains a large number of defects and fissures that determine the crack morphology after fracture under dynamic loading^[2–3]. Therefore, it is of great significance to study the dynamic characteristics and crack propagation laws of rock materials under impact loading.

Split Hopkinson pressure bar (SHPB) test system is widely used for dynamic impact loading test to analyze the response, crack propagation and energy dissipation characteristics of materials under impact loading^[4–7]. By using the improved single cleavage semi-circle (ISCSC) specimen configuration, Wang et al.^[8] studied the overall process of crack propagation and termination under impact loading. Cao et al.^[9] investigated the dynamic fracture of single cleavage drilled compression (SCDC) specimens, and recorded the times of dynamic initiation, propagation, termination and secondary initiation of cracks. Ni et al.^[10–11] studied the dynamic fracture toughness and stress intensity factors of SCDC specimens and double cleavage drilled compression (DCDC) plates under impact loading. In general, using specimens with pre-existing cracks to study the crack propagation under impact loading has become the dominant method in this field. Among them, the single cleavage triangle (SCT) specimen adopted by Dong et al.^[12–13] not only facilitates the study of the crack growth law, but also has the advantages of simple structure and easy processing, and it has also been adopted in the research of Wang et al.^[14–15].

Previous studies focused on the whole propagation process of mode I crack, especially the initiation, termination and secondary initiation. However, statistical studies on crack propagation patterns have been rarely reported, which can be quantitatively studied by the Hough transform method. This method was proposed by Hough^[16] and allows the detection of specific shapes in image space using the duality of line–point in the parameter space, including straight line^[17], circle^[18], ellipse^[19], etc. For the detection of crack images, David et al.^[20] developed

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the FracPaq toolbox based on Matlab software and conducted the relevant research. Rizzo et al.^[21–22] adopted the FracPaq toolbox to analyze the scanning electron microscope (SEM) crack images of sandstones, and used the two-dimensional (2D) continuous wavelet transform method to analyze the directional changes of cracks at different scales. In the subsequent studies, statistical functions such as maximum likelihood estimation and K-S test were added to the toolbox and then applied to the investigation of crack distribution of fractured rock mass in the Santa Cruz area, USA^[22].

In addition, the research on the fracture surface roughness of rocks after dynamic fracture has not attracted enough attention, although the quantitative characterization of structural surface roughness has been studied for about half a century^[23]. The main debates focus on the influence of roughness on the rock mass under shear effect^[24–26], while there are few studies on the influence of roughness on the rock mass under impact loading. Therefore, the fracture characteristics of SCT specimens under impact loading can be investigated with the aid of existing quantitative indices of rock fracture surface roughness.

During the acquisition of fracture surface roughness, according to different detection principles, the detection devices are classified into probe contact^[23], photogrammetric^[27] and optical types^[28]. Among them, the fracture surface roughness detection method based on the optical principle possesses the advantages of high-precision and non-damage to the structural plane^[23], with the measurement accuracy even reaching nanometer level. However, in the point cloud data of fracture surface obtained by optical devices, some obvious error points are inevitable, which brings inconvenience to the subsequent processing. The reasons for these error points are as follows: (1) In the case of high precision, there is an upper limit of the measurement range in the vertical direction of the detection device. For specimens with excessive roughness, the recorded data of individual points may exceed this range, resulting in a large gap between the recorded value and the actual value. (2) The high transparency of the individual particles in the rock specimen can easily cause the scattering of the detector light source, making the detector unable to capture the reflected light source. This phenomenon

has not attracted enough attention yet.

To address the above problems, in this study, the impact test on the SCT granite specimens is carried out using the SHPB test system. The Hough transform method is adopted to quantitatively describe the morphology of mode I crack. The relationship between the impact pressure and the statistical results of the distributions of crack lengths and angles is expounded, and the fracture surface reconstruction method of fitting surface inversion is proposed. The problem of error points in the three-dimensional (3D) point cloud data of fracture surface is solved, and the relationship between rock fracture surface roughness and air pressure is discussed, which provides a valuable reference for the study of rock dynamic fracture.

2 Test system and principle

2.1 Specimen preparation

The granite specimens used in the test were sampled from the Dagushan Iron Mine in Anshan, China. All the specimens were taken from an intact rock mass with a good homogeneity. The physico-mechanical parameters of rock specimens are listed in Table 1. The SCT rock sample proposed by Dong et al.^[12] was used to prepare the specimens in a rectangular configuration with a triangular opening. The prefabricated crack was located at the apex of the triangle. The processed granite specimen is 50 mm wide, 100 mm high and 30 mm thick. A prefabricated crack of 10 mm in length was cut by a 1 mm thick saw blade and then sharpened. The pad is a triangular body made of steel (density of 7.9 g /cm³), with a top angle of 30°, a bottom length of 16.0 mm, a height of 29.6 mm and a thickness of 30.0 mm. The structural size of the SCT specimen is illustrated in Fig. 1.

2.2 Test device and schemes

The SHPB test system at the University of Science and Technology Liaoning was adopted in this test, as shown in Fig. 2. The lengths of the incident bar, transmitted bar and absorbing bar are 2 100 mm, 1 800 mm and 800 mm, respectively. Both the striker and the weight bar are made of high-strength steel with a diameter of 50 mm and an elastic modulus of 210 GPa. The 3D roughness detection was performed using the PS50 profilometer produced

 Table 1 Physico-mechanical parameters of rock specimens

Specimen ID	Diameter /mm	Height /mm	Density /(g • cm ⁻³)	Compressive strength /MPa	Elastic modulus /GPa	Deformation modulus /GPa	Poisson's ratio
1	50.08	100.65	2.64	168.46	46.25	32.23	0.28
2	50.06	100.44	2.65	156.57	43.95	20.76	0.27
3	49.41	100.60	2.64	147.32	42.55	20.18	0.14
4	50.03	100.47	2.65	159.40	41.06	23.12	0.28
5	50.06	100.49	2.64	159.69	39.71	16.25	0.31
Mean value	49.93	100.53	2.64	158.29	42.70	22.50	0.26

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(c) Side view



by NANOWEA company in the USA. The horizontal automatic scanning range is 50 mm, the scanning step size is $0.1 \mu m$, and the scanning speed is 20 mm/s.

During testing, the specimen embedded with the pad is placed between the incident bar and the transmitted bar. The side of the pad is connected with the incident bar, and the side of the specimen is connected with the transmitted bar. The air pressures in the test were set to 0.2, 0.3, 0.4, 0.5 and 0.6 MPa. After the impact fracture of the specimen, the fracture surface of the specimen was scanned by a 3D topography detection equipment to generate the 3D point cloud data of fracture surface.

Before testing, a typical test was performed to check the stress equilibrium. Figure 3(a) shows the typical stress wave pattern of a tested rock specimen, and the incident wave, transmitted wave and reflected wave can meet the test requirements. Figure 3(b) illustrates the typical stress equilibrium test curve, indicating the requirements of stress equilibrium during dynamic loading.







Fig. 3 Trial experiment and dynamic stress equilibrium

2.3 Calculation of energy dissipation

In the SHPB test, the spindle-shaped striker is launched by the high-pressure gas and gains the kinetic energy to impact the incident bar. The kinetic energy of the striker is converted into incident energy. The incident bar impacts the rock specimen, in which process the incident energy is converted into reflected energy, absorbed energy of the specimen and transmitted energy through the specimen. By ignoring the kinetic energy of rock fragments and the heat exchange between the rock and outside, the energy composed of three parts can be calculated according to the law of conservation of energy^[29]:

$$W_{i} = \frac{A_{0}C_{0}}{E_{0}} \int \sigma_{i}^{2} dt = A_{0}E_{0}C_{0} \int \varepsilon_{i}^{2} dt$$

$$W_{r} = \frac{A_{0}C_{0}}{E_{0}} \int \sigma_{r}^{2} dt = A_{0}E_{0}C_{b} \int \varepsilon_{r}^{2} dt$$

$$W_{t} = \frac{A_{0}C_{0}}{E_{0}} \int \sigma_{t}^{2} dt = A_{0}E_{0}C_{0} \int \varepsilon_{t}^{2} dt$$

$$(1)$$

where W_i , W_r and W_t represent the incident energy, reflected energy and transmitted energy, respectively; A_0 and E_0 are the sectional area and elastic modulus of the bar, respectively; C_0 is the longitudinal wave velocity; σ_i , σ_r and σ_t correspond to the stresses in the incident bar, reflected bar and transmitted bar, respectively; ε_i , ε_r and ε_t represent the strains of the incident bar, reflected bar and transmitted bar, respectively; and *t* is the corresponding time.

The energy absorbed by the specimen W_d can be calculated by

$$W_{\rm d} = W_{\rm i} - W_{\rm r} - W_{\rm t} \tag{2}$$

The energy density index w_s is introduced to evaluate the absorption of energy in rock using following equation:

$$w_{\rm s} = \frac{W_{\rm d}}{V} = \frac{A_0 E_0 C_0}{A_{\rm s} l_{\rm s}} \int \left[\mathcal{E}_{\rm i}(t)^2 - \mathcal{E}_{\rm r}(t)^2 - \mathcal{E}_{\rm t}(t)^2 \right] \mathrm{d}t \qquad (3)$$

where V is the volume of the specimen; A_s is the sectional area of the specimen; and l_s is the length of the specimen.

On this basis, Wang et al.^[29] proposed the energy time density at peak point to evaluate the rock energy dissipation under impact loading, and its calculation formula is written as

$$w_{\rm td}(t) = \frac{w_{\rm s}(t_{\rm p})}{t_{\rm p}} \tag{4}$$

where t_p is the moment when the peak point of energy time density appears; $w_s(t_p)$ is the energy consumption density at the time t_p ; and $w_{td}(t)$ is the modified calculation method of energy time density.

3 Processing of point cloud data and crack image

3.1 Error points removal of point cloud data

When detecting the rock fracture surface roughness, the detection principle of the testing equipment is to extract and process the point cloud data of the fracture surface. Due to the limitation of equipment measuring range and

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Fig. 4 Influence of error point on fracture surface

3.2 Construction of fitting surface

When the point cloud data are used to construct a surface, a plane (x, y) is first constructed as the domain boundary, and each data point should be located in the matrix mesh of the surface in the domain. For a node at a certain position in the matrix mesh, there are generally three methods to calculate the value of this point: nearest interpolation method, trilinear interpolation, and bilinear interpolation. Based on these three methods, the value of each point can be predicted according to the local value in the mesh, so that the validity of each data point can be guaranteed. In the fitting process, it is generally acknowledged that the mesh is essentially composed of lower-order splines, and then the behavior of a given node in the mesh becomes the interpolation problem of a linear combination, which can be described by linear algebra as follows:

$$Ax = y \tag{5}$$

where the length of the vector \mathbf{x} is equal to $n_x \cdot n_y$ (n_x and n_y are the numbers of the nodes in the x and y directions). Therefore, \mathbf{A} is a coefficient matrix with n rows, corresponding to each data point $n_x \cdot n_y$ in the fracture surface point cloud.

The essence of surface fitting is to deal with a regularization problem, which makes the (first) partial derivatives of the surface in the adjacent elements equal at the data points, and then produces the second set of linear equations described in the following form:

$$Bx = 0 \tag{6}$$

where the derivative is approximated by the finite difference method or Laplace transform method on the surface on adjacent nodes. First, the matrices A and B are scaled so that each matrix has a norm of 1, and then the vector xis solved to minimize the following equation:

$$\| (\mathbf{A}\mathbf{x} - \mathbf{y}) \|^2 + \lambda \| \mathbf{B}\mathbf{x} \|^2$$
(7)

where λ is the adjustment coefficient and its value is 1. Therefore, in the fitting process, it is necessary to set the numbers of nodes n_x and n_y required for fitting in the x and y axes in the independent mesh. The number of points in the x and y directions affects the number of meshes in the fitting surface. The correlation between the number of points and the number of meshes obtained is shown in Fig. 5(a). In the positive direction of both axes, the number of meshes increases with the increase in number of points. Theoretically, the finer the mesh on the surface, the better the reproduction of the fracture surface, and the validity of all point cloud data can therefore be guaranteed; otherwise, the triangle distortion shown in Fig. 4 will occur.



Fig. 5 Relationship between number of points and number of grids in x and y axes

The fracture surface roughness of the specimen under the impact load of 0.2 MPa was identified, and three sets of data with different node numbers on the x and y axes were selected for surface fitting, as presented in Fig. 5(b), 5(c) and 5(d). In Fig. 5(b), the numbers of nodes in the two directions are 300 and 200. The number of meshes has reached 6 000, and the meshes are relatively dense. However, it has been found that the error data have an impact on the results. The mesh shape at the edge of the fracture surface is irregular, indicating that the fitting surface results are greatly affected by the error point data, and the current selection of the node number is unreasonable. In Fig. 5(d), the numbers of nodes in the two directions are 4 and 30, and the number of meshes decreases to 120. Such a scarce mesh number suggests that the fitting surface can only represent the macro-level trend of the fracture surface, thereby losing a great deal of the details of rock fracture surface. Obviously, the setting of the value in Fig. 5(d) is far from the research goal. Among the three groups of tests, the more reasonable values are the results with the numbers of nodes being 8 and 150 and the number of meshes being 1 200. It can be clearly seen from Fig. 5(c) that the results of the fitting surface can effectively avoid the influence of error points without losing details. In the original point cloud data of fracture surface under the impact load of 0.2 MPa, there are 3 000 data points in the x direction and 200 data points in the y direction. Therefore, the reasonable range for the number of nodes in the two directions is 0.03-0.08 of the total number of point clouds in the corresponding direction.

3.3 Reconstruction of fracture surface data

According to the above method, the fitting surface is obtained, as shown in Fig. 6(a). The number of nodes in the x and y directions is set to 0.03-0.08 of the total number of points in the corresponding direction. Because the number of points on the fitting surface is much less than the actual point cloud data, it is inconvenient to compare the fitting surface with the original point cloud data. Before comparing the two groups of matrices, the matrices should be expanded by referring to the fitting results, and the data difference method was adopted to keep the numbers of rows and columns of the surface the same as that of the original point cloud data matrix, as shown in Fig. 6(b). The reconstructed point cloud data are all located on the fitted surface. The difference between the expanded matrix and the original point cloud data was calculated. The distribution of the absolute values of the difference at each position in the matrix is shown in Fig. 6(d), and it

physically measures the distance between the original data and the fitting surface. In the further processing of point cloud data, the threshold can be set according to the absolute value of the difference between the two matrices. The deletion principle is that when the absolute difference is greater than the threshold, the data at this point will be deleted and replaced with the data point on the fitting surface.

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Figure 7(a) shows the original distribution of the

thresholds. The statistical histogram of absolute differences is shown in Fig. 7(b). The proportion of absolute differences less than 0.5 is as high as 80%, indicating a small proportion of error points in the original data. In Fig. 7(c), the distribution of error points inside the rock fracture surface is relatively scattered, and the high absolute differences are concentrated at the edge, indicating a high probability of error points appearing at the edge of the fracture surface. Figure 7(d) shows the relationship between



Fig. 7 Threshold setting

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the thresholds and the remaining points in the point cloud data. In terms of the overall trend, the remaining points increase significantly at the initial stage when the values are less than 1, and most of the error values are within the threshold range from 0 to 1. When the threshold is greater than 2, the change of the remaining points is approximately similar to the line parallel to the x axis, and the proportion of the remaining points exceeds 80%. A power function is used to fit the data points as follows:

$$y = -562184.1e^{\left(-\frac{x}{0.12}\right)} + 565188.16, \ R^2 = 0.93$$
 (8)

A closer look at the threshold within 0-1 reveals that the fitting curve tends to be parallel to the *x* axis when the threshold is 0.25. Admittedly, there are some errors in the interpolation of the fitting surface. Considering that the purpose of filtering is to replace large absolute difference points, a certain error value is acceptable. Therefore, the absolute difference of 0.25 can be selected as the threshold to achieve the removal of error points. **3.4 Crack image processing method**

The Hough transform method proposed by Paul Hough is adopted to process the crack propagation morphology, as shown in Fig. 8(a). After obtaining the image of crack region in the rock specimen, the image is binarized for processing. Then, the binarized image is processed by Hough transform for quantitative crack statistics.

As shown in Fig. 8(b), in the test, each fractured specimen is divided into left and right regions, i.e. region A and region B. The cracks corresponding to the front and back sides are termed A1, A2 and B1, B2, i.e. A1 and A2 coalesce in region A, and B1 and B2 coalesce in region B. The angles between the central axis and the cracks A1, A2, B1 and B2 are θ_{A1} , θ_{A2} , θ_{B1} and θ_{B2} , respectively.

4 Analyses of crack morphology and fracture surface

4.1 Effect of impact load on crack length distribution

The general rule of the test results is that the main crack of each group initiates from the base of the prefabricated crack, and propagates in the transverse direction rather than the vertical direction, which is somewhat different from the observations by Dong et al.^[13] and Wang et al.^[8]. In the study of Wang et al.^[8], tight sandstone was selected as the test material, and it was found that the density of rock materials and the composition of large particles affect the crack propagation behavior. The difference between our test results and the previous ones is that the heterogeneity of the rock leads to the deviation of the crack propagation path.

The crack images and maximum likelihood statistical results are shown in Fig. 9. After scale conversion, each



(a) Diagram of the Hough transform



Fig. 8 Crack image processing method

pixel in the images is 0.1 mm in size. The results suggest that the larger the value of λ , the fewer the large values of x. The λ value of each group fluctuates between 0.018 and 0.030, indicating that the overall crack distributions under different test conditions are similar. After observing the crack images of the front and back sides, it is found that in the first 4 groups of tests, the difference of λ value is only about 0.01, indicating that the specimen has a certain degree of symmetry and fracture uniformity during the fracturing. In comparison, the difference of λ in group 2 test with energy time density of 0.005 42 J/($cm^3 \cdot \mu s$) is 0.04. The reason for this is that there is an obvious gap on the left side of the front of the specimen, while the gap is smaller on the back of the specimen. For the last group, the energy time density is 0.012 93 J /(cm³ • μ s), and the difference of λ is as high as 0.12. The difference in macrocrack morphology is not obvious, but the symmetry of the cracks on both sides is poor, and the fracture surface is heterogeneous, resulting in a large gap of λ in this group. In Fig. 9(f), λ gradually increases with the increase in energy time density, suggesting that the number of cracks in the specimen also tends to increase with the increase in absorbed energy.

Figure 10 shows the statistical distribution of crack





Fig. 9 Quantitatively statistical results of cracks and energy dissipation



Fig. 10 Statistics of crack length

https://rocksoilmech.researchcommons.org/journal/vol44/iss7/1 DOI: 10.16285/j.rsm.2022.6209 lengths in each group. In the overall trend, crack lengths ranging from 0 to 0.5 mm account for a large proportion, followed by cracks ranging from 0.5 mm to 1.0 mm, and the cracks larger than 1.5 mm share a small proportion. The results demonstrate a good crack connectivity. The crack width is narrow, and a few long straight lines can be combined to reconstruct a complete crack image. The high peak values of short cracks indicate the high degree of fracture of rock fracture surface. In addition, the peak frequency of the specimen on the back side appears in the range of 20.78–26.95 mm, close to the peak frequency appearing between 18.93-30.48 mm on the front side. The peak frequency is the lowest when the energy time density reaches 0.005 42 J /($cm^3 \cdot \mu s$) (air pressure of 0.3 MPa). The possible reason is that there is a notch at the edge of the specimen, and the crack extension is relatively long, while the fracture degree at the fracture surface is not high under low energy absorption, resulting in a minimum peak crack length after Hough transform. Another possible reason is that the impact pressure is low, so that the limited energy absorbed by the specimen leads to the generated fracture close to the prefabricated crack tip, giving rise to more short cracks under this pressure. Under other test conditions, the peak frequency is clustered at the crack length of 0.03 mm, while in the test with energy time density of 0.001 96 J /(cm³ \cdot µs) (air pressure of 0.2 MPa), the peak frequency is encountered

at 0.15 mm. Under low pressure conditions, the cracks mainly consist of short and medium cracks. Generally, the peak frequency increases with the increase in air pressure. The peak crack lengths are 2.52, 2.25, 2.79, 2.73 and 3.05 mm on the front side, and 2.52, 1.79, 2.89, 2.29 and 2.73 mm on the back side. The reason for the large data discreteness is that the crack propagation direction of each group is different to some extent, and the peak crack length estimated by Hough transform is not the real crack length, which has little reference significance and results in data discreteness.

4.2 Distribution of crack angles

The rose diagrams in Fig. 11 show the distribution of crack angles. The true north direction N is defined as the specimen vertical direction, the corresponding angle is 90°, and the horizontal direction is defined as 0°. The distribution of crack angles presents a central symmetry. The peak frequencies in the front side group are 11.41, 8.53, 7.87, 10.54 and 7.72, respectively, while those in



Fig. 11 Statistics of crack angle

the back side group are 8.40, 11.79, 8.22, 8.64 and 8.02, respectively. The mean value is 10%. The distribution of crack angles is relatively concentrated, and the frequency of most angles is between 5% and 10%. In contrast, the proportion of crack angle in the horizontal direction is the least, which indicates that although the cracks coalesce at the left and right sides of the specimen, the horizontal cracks are scarce. In the test results of the front and back sides, the cracks at 90° (270°) only account for a small proportion, illustrating that the vertical crack development is insignificant. The cracks develop more obvious at the angles of 15°-60° and 105°-150°, proving that the crack direction deviates from the vertical direction of the specimen at a certain angle. In addition, when the energy time densities are 0.012 22 J/($cm^3 \cdot \mu s$) (air pressure of 0.5 MPa) and $0.012 93 \text{ J}/(\text{cm}^3 \cdot \mu \text{s})$ (air pressure of 0.6 MPa), the angle distributions of the two groups of tests are relatively balanced, except for the vertical and horizontal directions. However, when the energy time densities are 0.001 96 J /(cm³ • μ s) (air pressure of 0.2 MPa), 0.005 42 J /(cm³ • μ s) (air pressure of 0.3 MPa) and 0.007 77 J /(cm³ • μ s) (air pressure of 0.4 MPa), the fluctuation of crack angles is more violent with some dominant directions appearing on the rose diagram of crack angle. It reveals that the directionality of cracks is less obvious and the bending degree of crack is larger in the case of high energy dissipation, whereas under low energy dissipation, the directional consistency of cracks is more prominent.

Previous studies^[10–12, 27] have shown that the growth process of prefabricated crack in the rock specimens under impact loading can be classified into three stages: cracks initiation, propagation and termination. In addition, the impact disturbance may lead to secondary initiation of cracks. In the results of five groups of tests, the direction of crack growth is significantly altered under the air pressure of 0.2-0.4 MPa, and the most obvious change occurs under the energy time density of 0.001 96 J/($cm^3 \cdot \mu s$) (air pressure of 0.2 MPa). As shown in Fig. 10(f), the absorbed energy increases linearly with the increase in air pressure. The energy absorbed by the SCT specimen is relatively small when the energy time density remains low. During the crack propagation, the phenomena of crack termination and secondary initiation occur, and the changing position is not far from the edge of the specimen. In addition, due to the heterogeneity of the granite in the preparation of the SCT specimen, the crack direction changes during secondary initiation of cracks, which is mainly caused by high energy dissipation. Therefore, in the test results of two groups of energy time densities of 0.007 77 J /($cm^3 \cdot \mu s$) (air pressure of 0.4 MPa) and $0.012 22 \text{ J/(cm}^3 \cdot \mu s))$ (air pressure of 0.5 MPa), the crack

termination and secondary initiation are rarely observed during the crack propagation along the main crack direction until coalescence.

In Figs. 8 and 11, although the cracks on the front and back sides of the specimen are symmetrical, the differences between the cracks on both sides are obvious, which are well demonstrated in the crack images and statistical results. In this regard, the absolute differences of crack length and angle are defined, representing the differences of crack length and angle between the front and back sides.

The absolute value of the crack length is calculated as $S_{1} = |L_{A1} - L_{A2}| + |L_{B1} - L_{B2}|$ (9)

Table 2	Statistics	of crack	morphology
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where S_1 is the sum of the absolute differences of crack length; *L* is the crack length; and the subscripts A1, A2, B1 and B2 correspond to different crack numbers.

The absolute value of the crack angle can be obtained by

$$S_{\theta} = \left| \theta_{\rm A1} - \theta_{\rm A2} \right| + \left| \theta_{\rm B1} - \theta_{\rm B2} \right| \tag{10}$$

where S_{θ} is the sum of the absolute differences of the crack angle; and θ is the crack angle.

The statistics of crack length, crack angle and their absolute values are listed in Table 2, and the relationship between energy time density and absolute differences is plotted in Fig. 12.

Air pressure /MPa	Energy time density $/(J \cdot cm^{-3} \cdot \mu s^{-1})$	Crack length at front side /mm		Crack length at back side /mm		Crack angle at front side /(°)		Crack angle at back side /mm		Absolute difference	Absolute difference
		A1	B1	A2	B2	Al	B1	A2	B2	in lengui /min	In angle /()
0.2	0.001 96	23.1	26.2	24.1	24.8	60	45	75	50	2.4	20
0.3	0.005 42	41.9	26.5	37.4	21.0	50	90	45	75	10.0	20
0.4	0.007 77	34.1	25.2	36.5	20.2	55	55	45	60	7.4	15
0.5	0.012 22	29.4	28.0	31.5	21.0	55	75	45	75	9.1	10
0.6	0.012 93	25.8	28.2	33.0	31.0	55	45	55	40	10.0	5

As shown in Fig. 12, the absolute differences of crack length tend to increase with the increase in air pressure, but in the test with energy time density of 0.005 42 J /(cm³ • μ s) (air pressure of 0.3 MPa), the absolute differences are discrete, which may be caused by the notch on the front side of the specimen. It suggests that the more energy absorbed by the specimen, the more obvious the difference of cracks on both sides of the specimen. The absolute difference of crack angle of each group is close to each other, falling between 5° and 17°. In the test with energy time density of 0.012 93 J/($cm^3 \cdot \mu s$) (air pressure of 0.6 MPa), the absolute difference of crack angle is only 5°, indicating that the cracks on both sides of the specimen are closer to parallel under the test conditions. The results also confirm a low roughness of the fracture surface.



Fig. 12 Relationship between energy time density and crack statistical results

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4.3 Statistical results

The statistical characteristics of a 3D fracture surface can be defined according to the indicators proposed by Belem et al.^[25]:

$$Z_{2s} = \left\{ \frac{1}{L_x L_y} \int_0^{L_x} \int_0^{L_y} \left\{ \left[\frac{\partial z(x, y)}{\partial x} \right]^2 + \left[\frac{\partial z(x, y)}{\partial y} \right]^2 \right\} dx dy \right\}^{1/2}$$
(11)

Its approximation is expressed as

$$Z_{2s} = \left\{ \frac{1}{(N_x - 1)(N_y - 1)} \cdot \left[\frac{1}{\Delta x^2} \sum_{i=1}^{N_x - 1} \sum_{i=1}^{N_y - 1} \frac{(z_{i+1,j+1} - z_{i,j+1})^2 + (z_{i+1,j} - z_{i,j})^2}{2} + \frac{1}{\Delta y^2} \sum_{i=1}^{N_y - 1} \sum_{i=1}^{N_x - 1} \frac{(z_{i+1,j+1} - z_{i+1,j})^2 + (z_{i,j+1} - z_{i,j})^2}{2} \right] \right\}^{\frac{1}{2}} (12)$$

The 2D joint roughness coefficient (JRC) proposed by Barton et al.^[32] is positively linear with the statistical function. However, the above studies all took the shear test as the research object, and the dimensions of the selected samples were 100 mm×100 mm×100 mm. In this study, the specimens are all in long strip shape with dimensions of 100 mm×50 mm×30 mm. The fracture surface size is set to 20 mm×30 mm, and the sampling interval of 1 mm seems to be too large, resulting in Z_{2s} being too small. Figure 13 illustrates the variation of Z_{2s} of fracture surface A at different sampling intervals under energy time density of 0.001 96 J/(cm³ \cdot µs) (air pressure of 0.2 MPa).



Fig. 13 Relationship between sampling interval and Z_{2s}

With the increase in sampling interval, Z_{2s} decreases sharply until it approaches zero, showing an exponential decreasing trend. The peak Z_{2s} is 29.66, corresponding to $\Delta x = 0.01$ mm and $\Delta y = 0.1$ mm, which is the minimum step size of the 3D topography detection device in the x and y directions. By summarizing the Z_{2s} values of standard JRC curves in the literature, Yuan et al.^[24] found that they were all between 0 and 1, and concluded that excessive sampling interval is obviously unreasonable. For the convenience of calculation, the sampling interval of $\Delta x = 0.1$ mm and $\Delta y = 0.1$ mm is assumed, and the obtained Z_{2s} is 0.30. For the consistency of comparison, this sampling interval is also used in the calculation of other fracture surface statistical parameters.

The larger the value of Z_{2s} , the higher the roughness of rock fracture surface. As shown in Fig. 14, the correlation between the fracture surface roughness and air pressure is insignificant. However, the fracture surface roughness reaches the lowest under the air pressure of 0.6 MPa (energy time density of 0.012 93 J/(cm³ • µs)), and the highest under the air pressure of 0.3 MPa (0.005 42 J/(cm³ • µs)). The reason is that there is a noticeable notch on the left side of the specimen. The roughness parameter Z_{2s} of



Fig. 14 Relationship between air pressure and Z_{2s}

fracture surface A is 0.40, while for other groups, Z_{2s} values are similar. This is considered to be the outcome of crack termination and secondary initiation. In addition, the absolute differences of Z_{2s} for fracture surfaces A and B were calculated, and the roughness difference of each group follows a decreasing trend except for the test under the air pressure of 0.2 MPa (0.001 96 J /(cm³ • µs)). It shows that under high energy dissipation, the cracks in fracture surfaces A and B caused by impact loading coalesce well, the secondary initiation phenomenon is inconspicuous, and thus the difference in fracture surface roughness is negligible.

5 Conclusions

Based on the SHPB test system, this paper studied the dynamic fracture characteristics of SCT granite specimens, addressed the problem of removing error points in the 3D fracture surface point cloud data, and summarized the law of influence of impact load on the mode I crack morphology and fracture surface roughness. The main conclusions can be drawn as follow:

(1) The energy time density presents a positively linear relationship with the air pressure, and the functional relationship is $y = 0.028 \ 73x - 0.003 \ 43$ with $R^2 = 0.97$. With the increase in absorbed energy, the exponential distribution λ of crack length increases. Overall, the cracks with the length of 0–0.5 mm are the most, followed by the cracks with the length of 0.5–1.0 mm, and the proportion of cracks greater than 1.5 mm is the least.

(2) Morphologically, the cracks coalesce on the left and right sides of the specimen, and the horizontal cracks are less developed. The distribution of crack angles has no prominent direction, except at low energy time density, at which crack termination and secondary initiation would occur. At high energy time density, the cracks have the characteristics of higher bending degree, less obvious directionality, and good connectivity.

(3) The fitting surface threshold detection method is used to remove the error points and reconstruct the fracture surface data. It is found that the reasonable value of the fitting surface in the x and y directions is between 0.03 and 0.08, and the threshold is 0.25. The difference in roughness parameter of each group experiences a decreasing trend with the increase in energy time density, and the cracks of fracture surfaces A and B coalesce well at high energy time density.

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